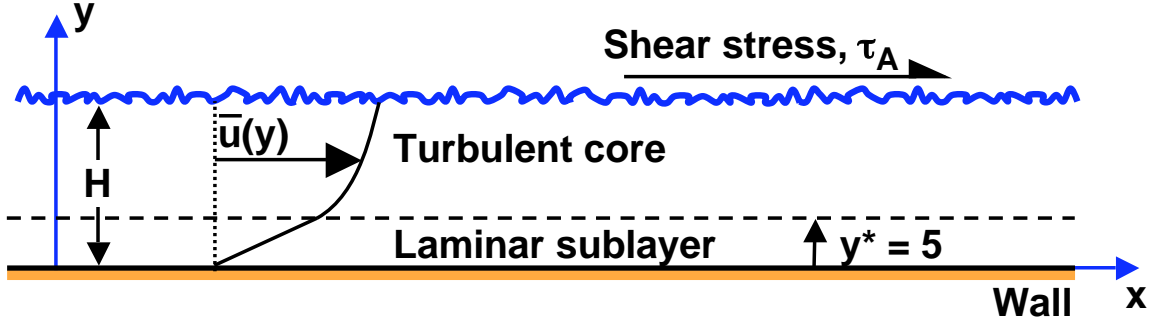


Solution to Problem 272A

Either by considering the momentum theorem applied to the control volume or by observing the similarity to Couette flow,



one can conclude that the shear stress, σ_{xy} , is a constant throughout the flow. Hence $\sigma_{xy} = \tau_A$.

- In the *turbulent core*, since the viscous stresses are negligible

$$\begin{aligned} \sigma_{xy} = \tau_A &= -\rho \overline{u'v'} \\ &= \rho \kappa^2 y^2 \left(\frac{d\bar{u}}{dy} \right)^2 \end{aligned}$$

by Prandtl's mixing length theory. Hence

$$\begin{aligned} \kappa y \frac{d\bar{u}}{dy} &= \sqrt{\frac{\tau_A}{\rho}} = u_\tau = \text{constant} \\ \frac{\bar{u}}{u_\tau} &= \frac{1}{\kappa} \ln y + C \end{aligned}$$

where C is an integration constant.

- Within the *laminar sublayer*, we now know $u_\tau = \sqrt{\frac{\tau_A}{\rho}}$

$$u^* = \frac{\bar{u}}{u_\tau} = y^* = \frac{u_\tau y}{\nu}$$

Therefore at the edge of the laminar sublayer

$$\begin{aligned} y^* &= 5 \\ \frac{\bar{u}}{u_\tau} &= 5 \\ y &= \frac{5\nu}{u_\tau} \end{aligned}$$

Hence to find the constant, C

$$C = 5 - \frac{1}{\kappa} \ln \left(\frac{5\nu}{u_\tau} \right)$$

Therefore in the turbulent core

$$\begin{aligned} \frac{\bar{u}}{u_\tau} &= \frac{1}{\kappa} \ln \left(\frac{u_\tau y}{5\nu} \right) + 5 \\ (\bar{u})_{y=H} &= \sqrt{\frac{\tau_A}{\rho}} \left[\frac{1}{\kappa} \ln \left(\frac{H}{5\nu} \sqrt{\frac{\tau_A}{\rho}} \right) + 5 \right] \end{aligned}$$