

## Flow over a Wavy Wall

A second example of this type of solution that was investigated by Zung (1967) is steady particle-laden

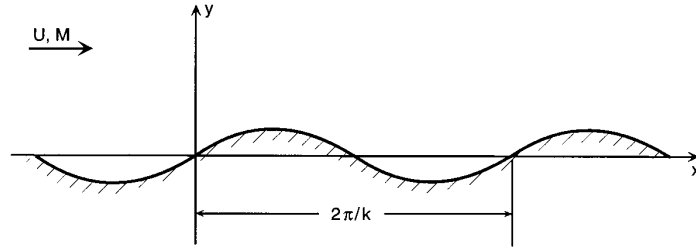


Figure 1: Schematic for flow over a wavy wall.

flow over a wavy wall of small amplitude (figure 1) so that only the terms that are linear in the amplitude need be retained. The solution takes the form

$$\exp(i\kappa_1 x - i\kappa_2 y) \quad (\text{Nnh1})$$

where  $2\pi/\kappa_1$  is the wavelength of the wall whose mean direction corresponds with the  $x$  axis and  $\kappa_2$  is a complex number whose real part determines the inclination of the characteristics or Mach waves and whose imaginary part determines the attenuation with distance from the wall. The value of  $\kappa_2$  is obtained in the solution from a dispersion relation that has many similarities to equation (Nnf8). Typical computations of  $\kappa_2$  are presented in figure 2. The asymptotic values for large  $t_u$  that occur on the right in this figure correspond to cases in which the particle motion is constant and unaffected by the waves. Consequently, in subsonic flows ( $M = U/c_C < 1$ ) in which there are no characteristics, the value of  $Re\{\kappa_2/\kappa_1\}$  asymptotes to zero and the waves decay with distance from the wall such that  $Im\{\kappa_2/\kappa_1\}$  tends to  $(1 - M^2)^{\frac{1}{2}}$ . On the other hand in supersonic flows ( $M = U/c_C > 1$ )  $Re\{\kappa_2/\kappa_1\}$  asymptotes to the tangent of the Mach wave angle in the gas alone, namely  $(M^2 - 1)^{\frac{1}{2}}$ , and the decay along these characteristics is zero.

At the other extreme, the asymptotic values as  $t_u$  approaches zero correspond to the case of the effective gas whose properties are given in section (Nnc). Then the appropriate Mach number,  $M_0$ , is that based on the speed of sound in the effective gas (equation (Nnc8)). In the case of figure 2,  $M_0^2 = 2.4M^2$ . Consequently, in subsonic flows ( $M_0 < 1$ ), the real and imaginary parts of  $\kappa_2/\kappa_1$  tend to zero and  $(1 - M_0^2)^{\frac{1}{2}}$  respectively as  $t_u$  tends to zero. In supersonic flows ( $M_0 > 1$ ), they tend to  $(M_0^2 - 1)^{\frac{1}{2}}$  and zero respectively.

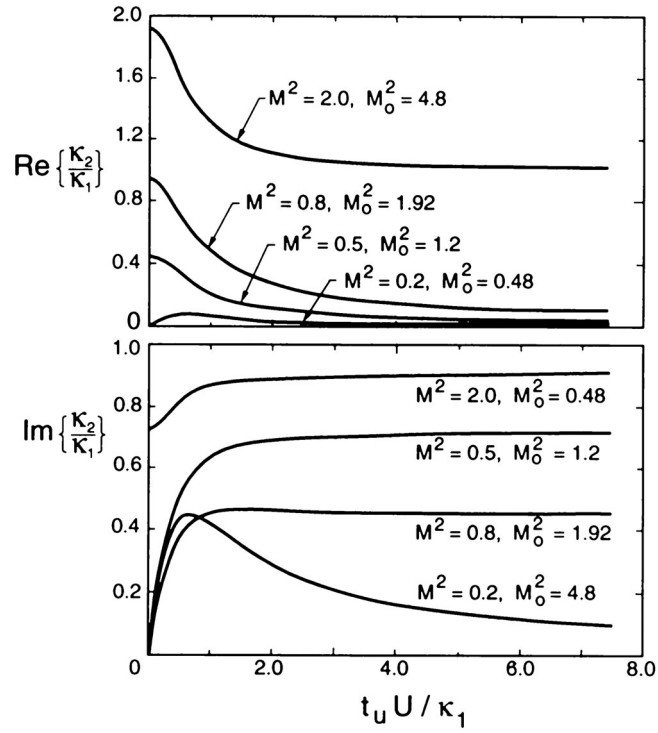


Figure 2: Typical results from the wavy wall solution of Zung (1969). Real and imaginary parts of  $\kappa_2/\kappa_1$  are plotted against  $t_u U/\kappa_1$  for various mean Mach numbers,  $M = U/c_C$ , for the case of  $t_T/t_u = 1$ ,  $c_{pC}/c_{sD} = 1$ ,  $\gamma = 1.4$  and a particle loading,  $\xi = 1$ .